

MC Simulation of the
Constrained Markovian Evolution
Loops and Legs 2004

S. Jadach

and

M. Skrzypek

`stanislaw.jadach@ifj.edu.pl, maciej.skrzypek@ifj.edu.pl`

IFJ-PAN, Kraków, Poland

The name of the game

$$(Loops + Legs)^\infty$$

The long standing problem

- Markovian MC implementing the QCD/QED evolution equations is basic ingredient in all parton shower type MCs

The long standing problem

- Markovian MC implementing the QCD/QED evolution equations is basic ingredient in all parton shower type MCs
- Unconstrained Markovian, with evolution kernels from perturbative QED/QED, can only be used for FSR (inefficient for ISR)

The long standing problem

- Markovian MC implementing the QCD/QED evolution equations is basic ingredient in all parton shower type MCs
- Unconstrained Markovian, with evolution kernels from perturbative QED/QED, can only be used for FSR (inefficient for ISR)
- For ISR the Backward Markovian of Sjostrand (Phys.Lett. 157B, 1985) is a widely adopted remedy.

The long standing problem

- Markovian MC implementing the QCD/QED evolution equations is basic ingredient in all parton shower type MCs
- Unconstrained Markovian, with evolution kernels from perturbative QED/QED, can only be used for FSR (inefficient for ISR)
- For ISR the Backward Markovian of Sjostrand (Phys.Lett. 157B, 1985) is a widely adopted remedy.
- **Backward Markovian does not solve evolution eqs. It merely exploits their solutions coming from the external non-MC methods**

The long standing problem

- Markovian MC implementing the QCD/QED evolution equations is basic ingredient in all parton shower type MCs
- Unconstrained Markovian, with evolution kernels from perturbative QED/QED, can only be used for FSR (inefficient for ISR)
- For ISR the Backward Markovian of Sjostrand (Phys.Lett. 157B, 1985) is a widely adopted remedy.
- Backward Markovian does not solve evolution eqs. It merely exploits their solutions coming from the external non-MC methods
- **Is it possible to invent an efficient MC algorithm for constrained Markovian based on *internal* MC solutions of the evolution eqs?**

Constrained Solutions are coming

- We have found a class of solutions of the above long-standing problem

Constrained Solutions are coming

- We have found a class of solutions of the above long-standing problem
- Introductory exercise: Markovian MC $E_{\text{VOL}}\text{MC}$ was found to agree with $\text{QCD}_{\text{num}}16$ to within 0.2%,
[Acta Phys.Polon. B35 \(2004\) 745](#)

Multicomponent evolution equation

$$\frac{\partial}{\partial t} D_k(t, x) = \sum_j \int_x^1 \frac{dz}{z} P_{kj}(z) \frac{\alpha_S(t, z)}{\pi} D_j\left(t, \frac{x}{z}\right)$$

Indices i and k denote gluon or quark,
Evolution time is $t = \ln(Q)$.

Multicomponent evolution equation

$$\begin{aligned}\frac{\partial}{\partial t} D_k(t, x) &= \sum_j \int_x^1 \frac{dz}{z} P_{kj}(z) \frac{\alpha_S(t, z)}{\pi} D_j\left(t, \frac{x}{z}\right) \\ &= \sum_j \frac{\alpha_S(t, \cdot)}{\pi} P_{kj}(\cdot) \otimes D_j(t, \cdot)\end{aligned}$$

$$f(\cdot) \otimes g(\cdot)(x) \equiv \int dx_1 dx_2 \delta(x - x_1 x_2) f(x_1) g(x_2)$$

Indices i and k denote gluon or quark,
Evolution time is $t = \ln(Q)$.

Multicomponent evolution equation

$$\begin{aligned}
 \frac{\partial}{\partial t} D_k(t, x) &= \sum_j \int_x^1 \frac{dz}{z} P_{kj}(z) \frac{\alpha_S(t, z)}{\pi} D_j\left(t, \frac{x}{z}\right) \\
 &= \sum_j \frac{\alpha_S(t, \cdot)}{\pi} P_{kj}(\cdot) \otimes D_j(t, \cdot) \\
 &= \sum_j \mathcal{P}_{kj}(t, \cdot) \otimes D_j(t, \cdot)
 \end{aligned}$$

$$f(\cdot) \otimes g(\cdot)(x) \equiv \int dx_1 dx_2 \delta(x - x_1 x_2) f(x_1) g(x_2)$$

$$\mathcal{P}_{kj}(t, z) \equiv \frac{\alpha_S(t, z)}{\pi} P_{kj}(z)$$

Indices i and k denote gluon or quark,

Evolution time is $t = \ln(Q)$.

Monte Carlo solution of Evolution Equation

$$\frac{\partial}{\partial t} D_k(t, x) = \sum_j \mathcal{P}_{kj}(t, \cdot) \otimes D_j(t, \cdot)$$

Differential equation \longrightarrow integral equation:

$$e^{\Phi_k(t, t_0)} D_k(t, x) = D_k(t_0, x) + \int_{t_0}^t dt_1 e^{-\Phi_k(t_1, t_0)} \sum_j \mathcal{P}_{kj}^\ominus(t_1, \cdot) \otimes D_j(t_1, \cdot)(x)$$

Monte Carlo solution of Evolution Equation

$$\frac{\partial}{\partial t} D_k(t, x) = \sum_j \mathcal{P}_{kj}(t, \cdot) \otimes D_j(t, \cdot)$$

Differential equation \longrightarrow integral equation:

$$e^{\Phi_k(t, t_0)} D_k(t, x) = D_k(t_0, x) + \int_{t_0}^t dt_1 e^{-\Phi_k(t_1, t_0)} \sum_j \mathcal{P}_{kj}^\ominus(t_1, \cdot) \otimes D_j(t_1, \cdot)(x)$$

where IR regulator is introduced

$$\mathcal{P}_{kj}(t, z) = -\mathcal{P}_{kk}^\delta(\epsilon(t)) \delta_{kj} \delta(1 - z) + \mathcal{P}_{kj}^\ominus(t, z) \Theta(1 - z - \epsilon)$$

Monte Carlo solution of Evolution Equation

$$\frac{\partial}{\partial t} D_k(t, x) = \sum_j \mathcal{P}_{kj}(t, \cdot) \otimes D_j(t, \cdot)$$

Differential equation \longrightarrow integral equation:

$$e^{\Phi_k(t, t_0)} D_k(t, x) = D_k(t_0, x) + \int_{t_0}^t dt_1 e^{-\Phi_k(t_1, t_0)} \sum_j \mathcal{P}_{kj}^\ominus(t_1, \cdot) \otimes D_j(t_1, \cdot)(x)$$

where IR regulator is introduced

$$\mathcal{P}_{kj}(t, z) = -\mathcal{P}_{kk}^\delta(\epsilon(t)) \delta_{kj} \delta(1 - z) + \mathcal{P}_{kj}^\ominus(t, z) \Theta(1 - z - \epsilon)$$

and the Sudakov formfactor appears

$$\Phi_k(t, t_0) = \int_{t_0}^t dt' \mathcal{P}_{kk}^\delta(\epsilon(t'))$$

Iterative multi-integral solution

$$\begin{aligned}
 D_K(t, x) &= e^{-\Phi_K(t, t_0)} D_K(t_0, x) \\
 &\times + \sum_{n=1}^{\infty} \sum_{K_0 \dots K_{n-1}} \prod_{i=1}^n \left[\int_{t_0}^t dt_i \Theta(t_i - t_{i-1}) \int_0^1 dz_i \right] \\
 &\times e^{-\Phi_K(t, t_n)} \int_0^1 dx_0 \prod_{i=1}^n \left[\mathcal{P}_{K_i K_{i-1}}^{\Theta}(t_i, z_i) e^{-\Phi_{K_{i-1}}(t_i, t_{i-1})} \right] \\
 &\times D_{K_0}(t_0, x_0) \delta(x - x_0 \prod_{i=1}^n z_i),
 \end{aligned}$$

where $K_n \equiv K$. Many options for the MC implementation.
 Generally they can be Markovian OR non-Markovian.

Iterative multi-integral solution

$$\begin{aligned}
 xD_K(t, x) &= e^{-\Phi_K(t, t_0)} D_K(t_0, x) \\
 &\times + \sum_{n=1}^{\infty} \sum_{K_0 \dots K_{n-1}} \prod_{i=1}^n \left[\int_{t_0}^t dt_i \Theta(t_i - t_{i-1}) \int_0^1 dz_i \right] \\
 &\times e^{-\Phi_K(t, t_n)} \int_0^1 dx_0 \prod_{i=1}^n \left[z_i \mathcal{P}_{K_i K_{i-1}}^{\Theta}(t_i, z_i) e^{-\Phi_{K_{i-1}}(t_i, t_{i-1})} \right] \\
 &\times x_0 D_{K_0}(t_0, x_0) \delta(x - x_0 \prod_{i=1}^n z_i),
 \end{aligned}$$

where $K_n \equiv K$. Many options for the MC implementation.

Generally they can be Markovian OR non-Markovian.

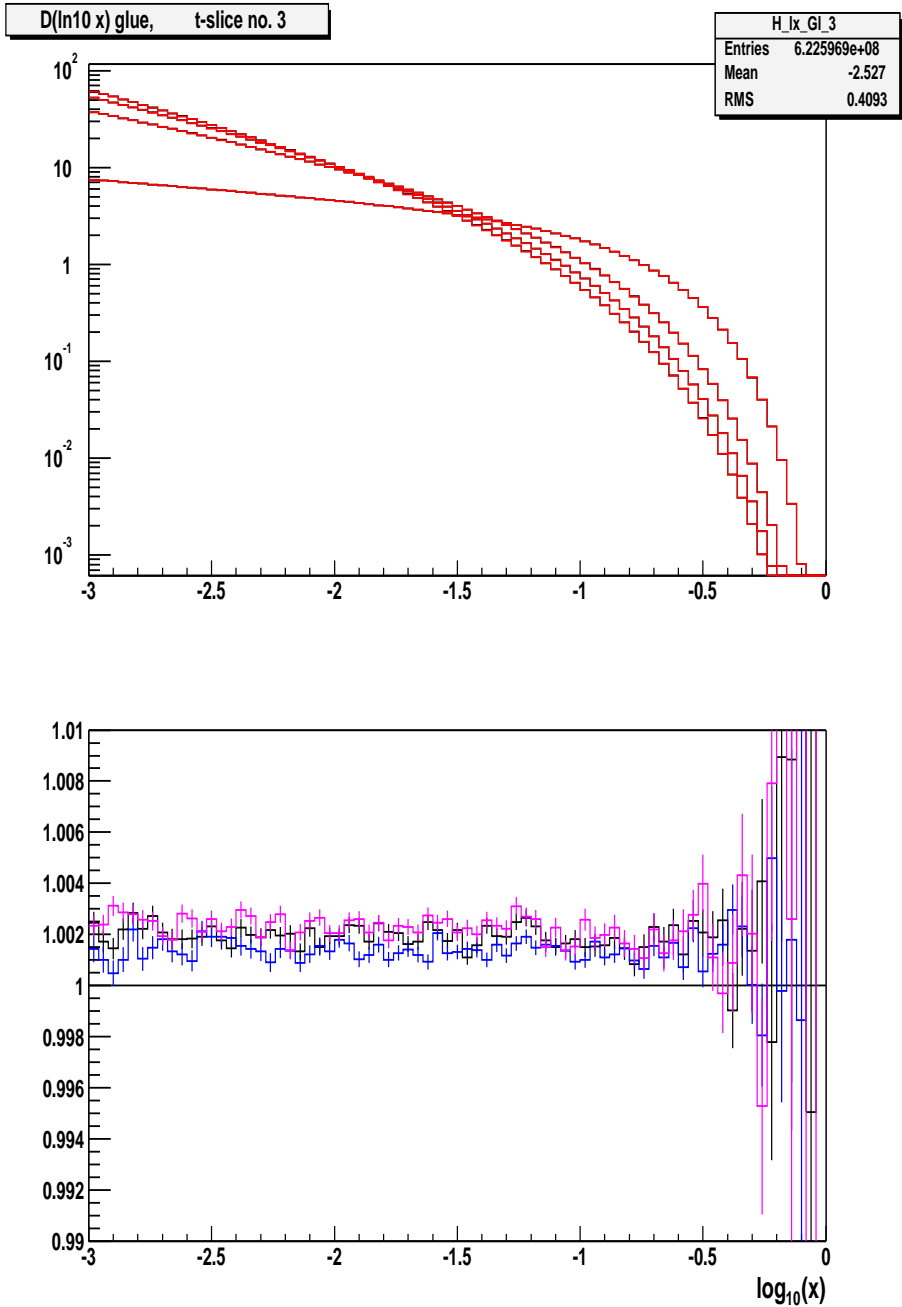
Solution for energy parton distributions $xD(x)$ more convenient!

Why? Kernels obey sum rules: $\sum_X \int dz z \mathcal{P}_{XK}(z) = 1$.

Master equation for Markovian solution

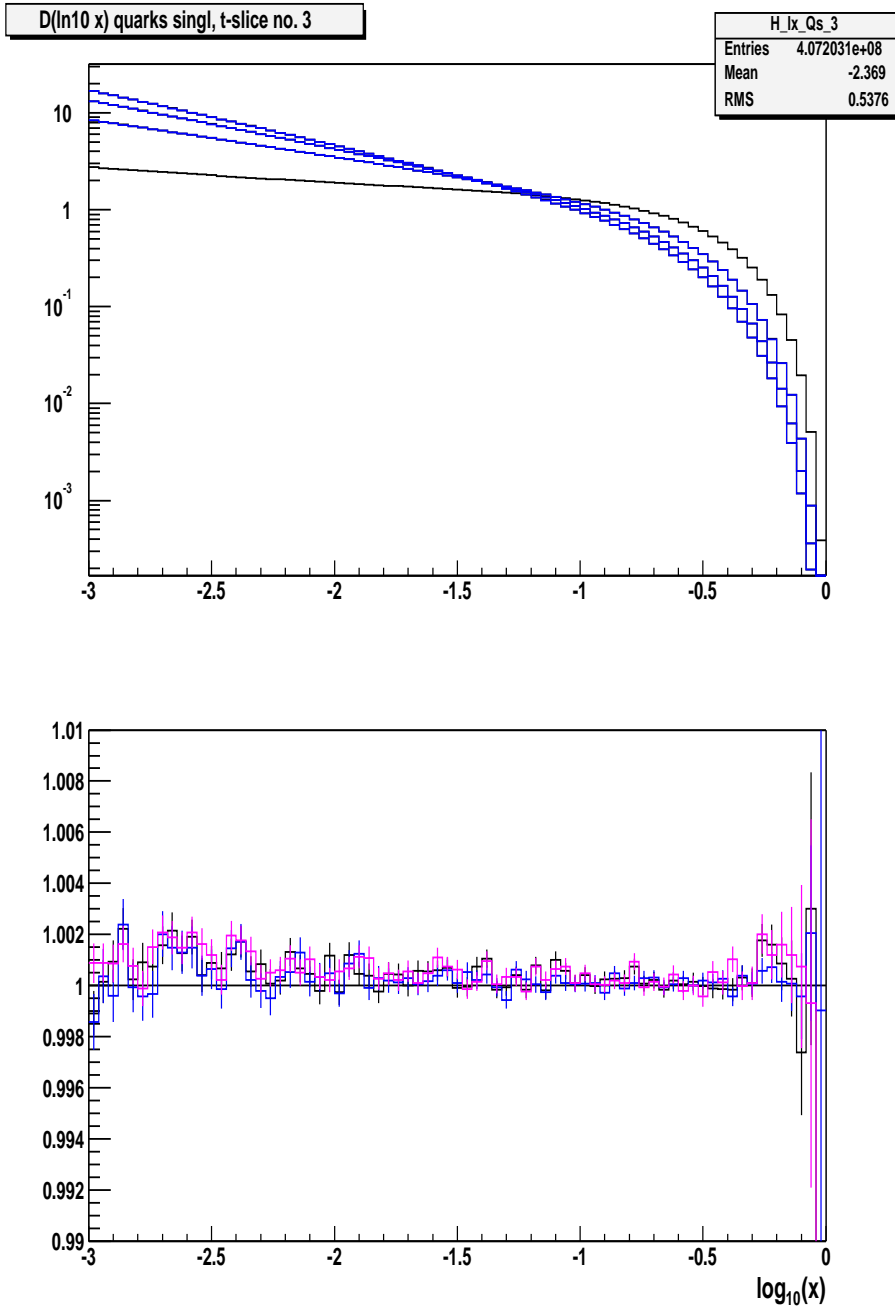
$$\begin{aligned}
 xD_K(\tau, x) &= \int_{\tau_1 > t} d\tau_1 dz_1 \sum_{K_1} \bar{\omega}(\tau_1, x_1, K_1 | \tau_0, x_0, K) xD_K(\tau_0, x) \\
 &+ \sum_{n=1}^{\infty} \int_0^1 dx_0 \int_{\tau_{n+1} > \tau} d\tau_{n+1} dz_{n+1} \sum_{K_{n+1}} \sum_{K_0 \dots K_{n-1}} \prod_{i=1}^n \int_{\tau_i < \tau}^t d\tau_i dz_i \\
 &\quad \times \bar{\omega}(\tau_{n+1}, x_{n+1}, K_{n+1} | \tau_n, x_n, K_n) \quad \leftarrow \text{spillover} \\
 &\quad \times \prod_{i=1}^n \bar{\omega}(\tau_i, x_i, K_i | \tau_{i-1}, x_{i-1}, K_{i-1}) \quad \leftarrow \text{normal step} \\
 &\quad \times \delta(x - x_0 \prod_{i=1}^n z_i) x_0 D_{K_0}(\tau_0, x_0) \bar{w}_P \bar{w}_\Delta \quad \leftarrow \text{MCweight}
 \end{aligned}$$

Tests: Proton \rightarrow gluon



Upper plot shows gluon distribution $x D_G(x, Q_i)$ evolved from $Q_0 = 1\text{GeV}$ to $Q_i = 10, 100, 100\text{GeV}$ obtained from QCDnum16 and EvolMC1, while lower plot shows their ratio. The horizontal axis is $\log_{10}(x)$. Starting distribution is complete proton at $Q = 1\text{GeV}$.

Tests: Proton \rightarrow quarks



Proton composition at 1 GeV

This is what we took for the introductory exercise:

$$xD_G(x) = 1.9083594473 \cdot x^{-0.2}(1-x)^{5.0},$$

$$xD_q(x) = 0.5 \cdot xD_{\text{sea}}(x) + xD_{2u}(x),$$

$$xD_{\bar{q}}(x) = 0.5 \cdot xD_{\text{sea}}(x) + xD_d(x),$$

$$xD_{\text{sea}}(x) = 0.6733449216 \cdot x^{-0.2}(1-x)^{7.0},$$

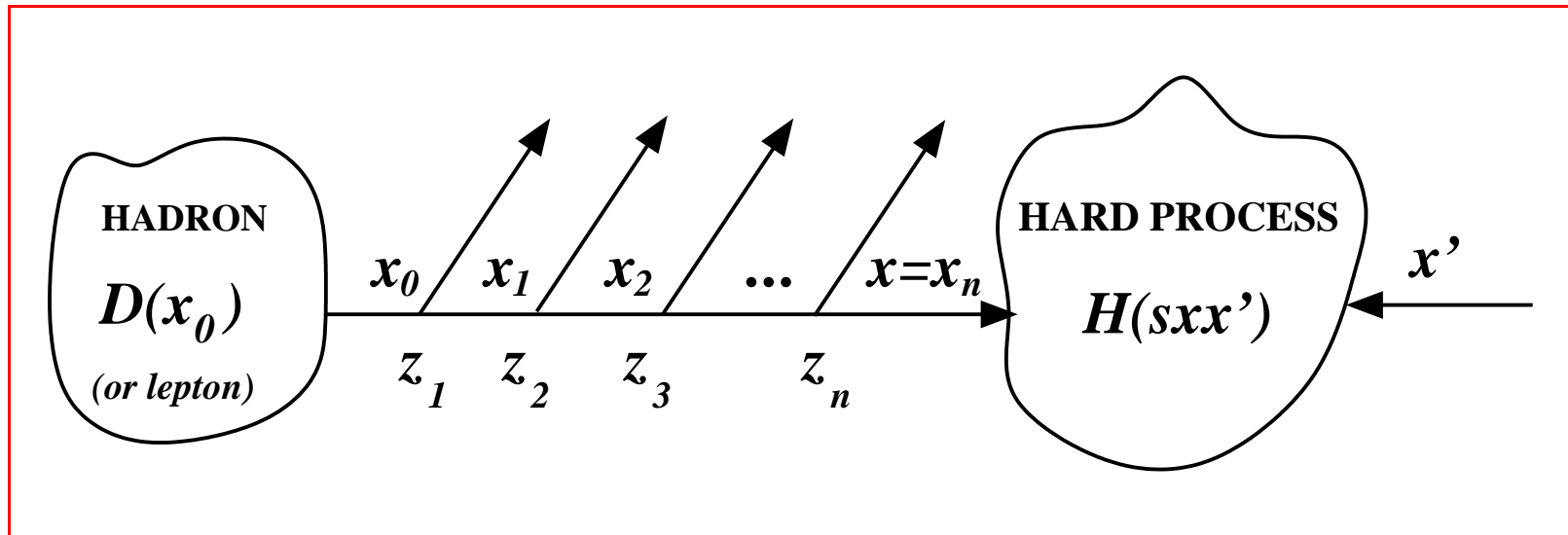
$$xD_{2u}(x) = 2.1875000000 \cdot x^{0.5}(1-x)^{3.0},$$

$$xD_d(x) = 1.2304687500 \cdot x^{0.5}(1-x)^{4.0},$$

Constrained Solutions are coming

- We have found a class of solutions of the above long-standing problem
- Introductory exercise: Markovian MC $E_{\text{vol}}MC$ was found to agree with $QCD_{\text{num}}16$ to within 0.2%,
[Acta Phys.Polon. B35 \(2004\) 745](#)
- **Recently, 1-st prototype of the efficient *constrained Markovian MC* (solution IIB) prototyped.**

Constrained Solutions class I and II



$$\int dx_0 D(x_0) \int \prod_i dz_i P(z_i) H(sx_0 \prod_i z_i)$$

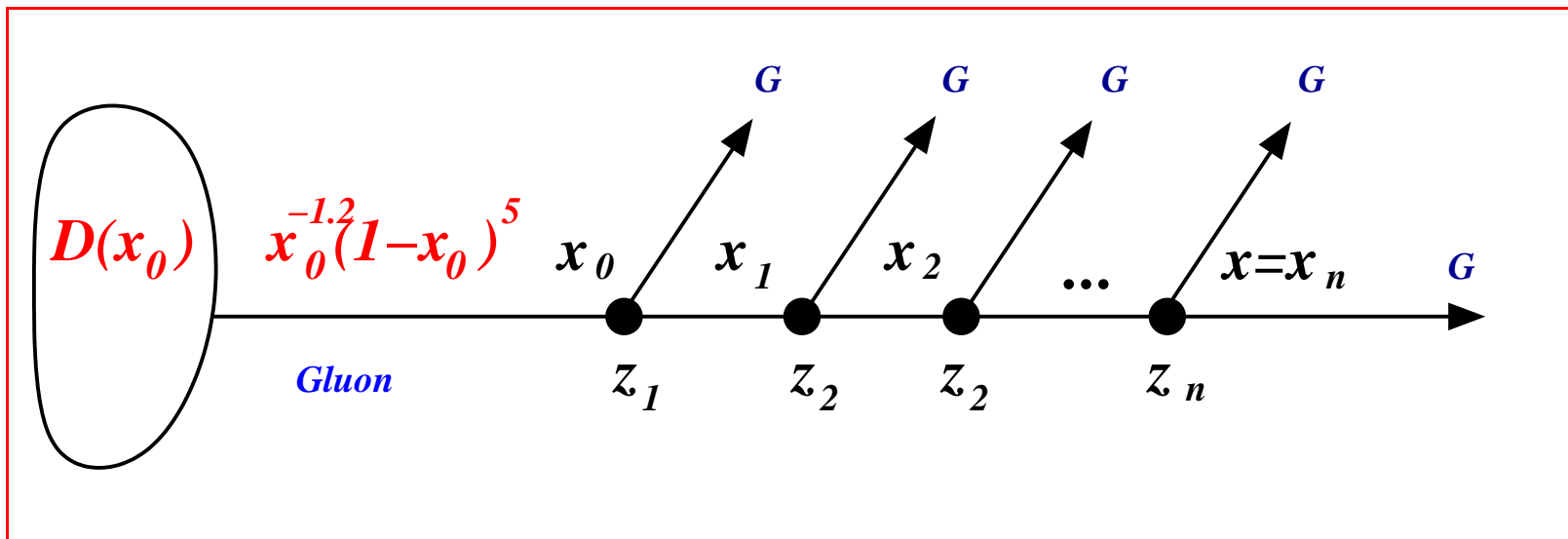
Solutions class I (more difficult because of $\delta(\dots)$):

$$\int dx dx_0 D(x_0) H(sx) \int \prod_i dz_i P(z_i) \delta(x - x_0 \prod_i z_i)$$

Solutions class II (only for QCD) **NEW!**:

$$\int dx H(sx) \int \prod_i \frac{dz_i}{z_i} P(z_i) D(x / \prod_i z_i) \Theta(\prod_i z_i - x)$$

Prototype IIB

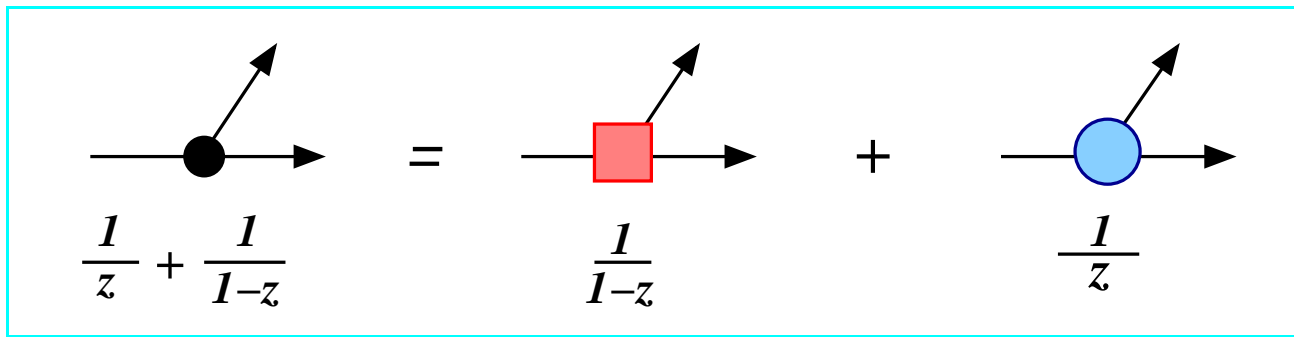


Replace $D(x_0) \rightarrow 1/x_0 = x \prod \frac{1}{z_i}$. Compensated by MC weight.

Must generate $P(z_i) = 2C_A \left(\frac{1}{z_i} + \frac{1}{1-z_i} \right)$

with the constraint $\prod_i z_i \geq x$. Not so trivial!

Solution by the multibranching method:

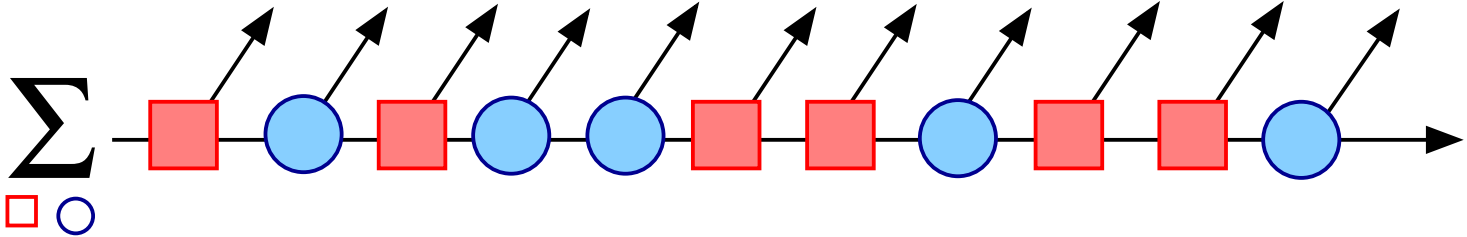


Multibranching in IIB

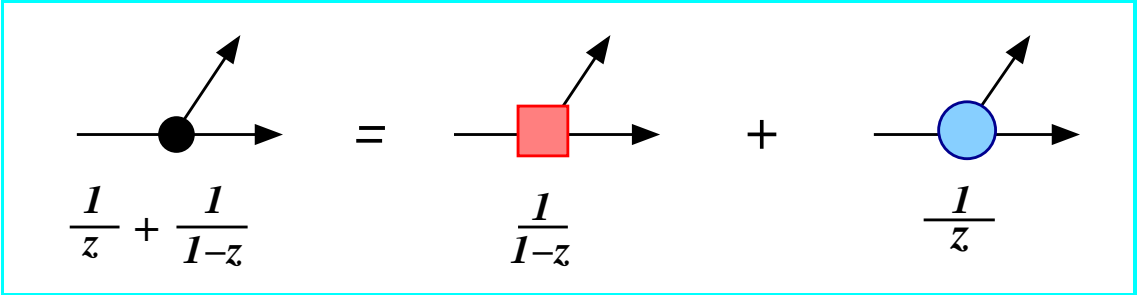
$$\frac{1}{z} + \frac{1}{1-z} = \frac{1}{1-z} + \frac{1}{z}$$

Using

Leads to sum over branches:

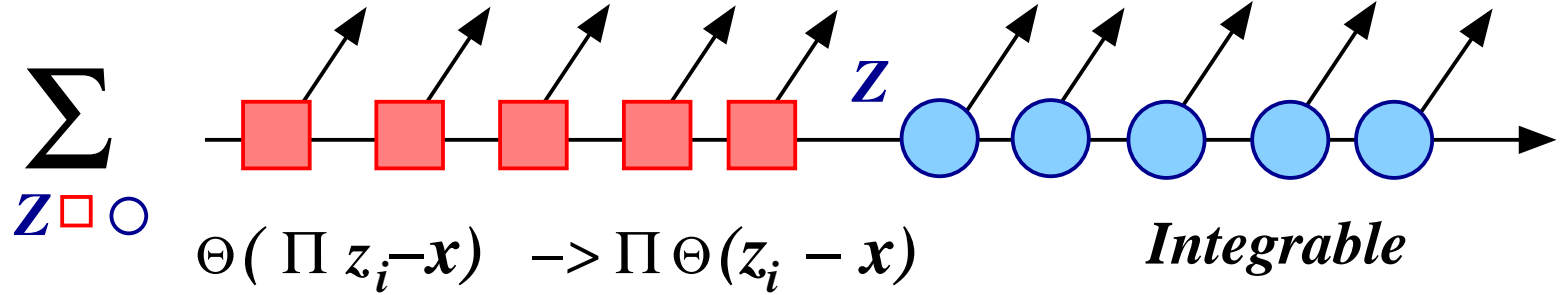


Multibranching in IIB



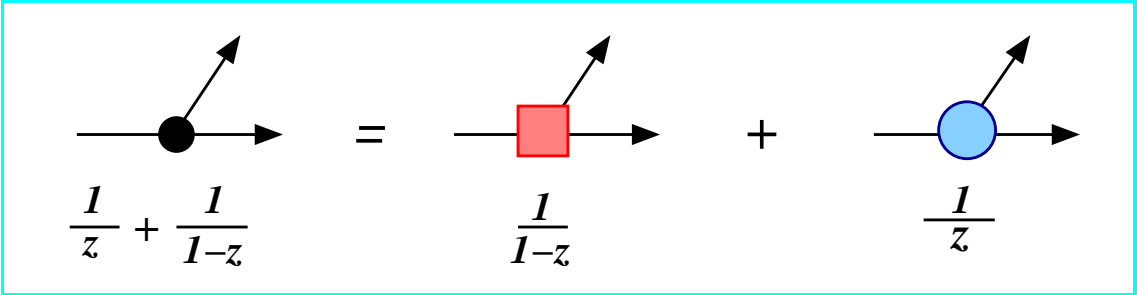
Using

Leads to sum over branches:



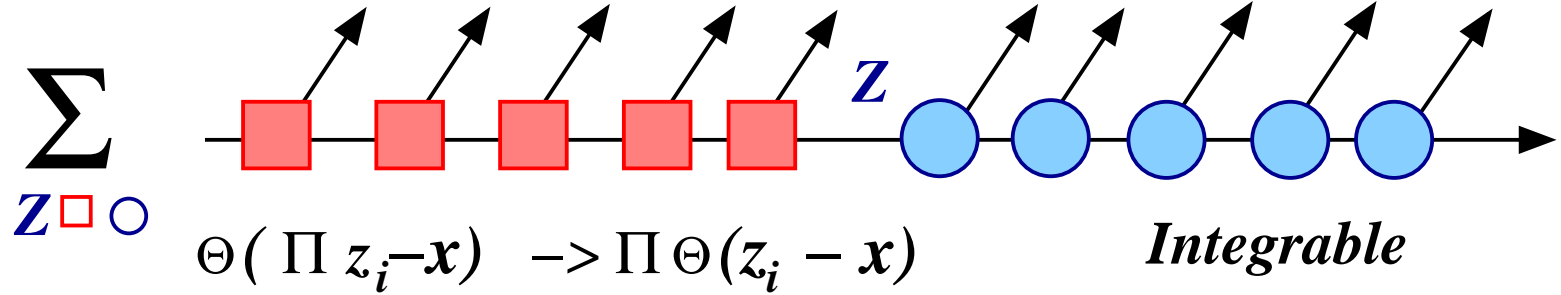
Contributions $1/z$ and $1/(1 - z)$ are combined and resummed separately.

Multibranching in IIB



Using

Leads to sum over branches:

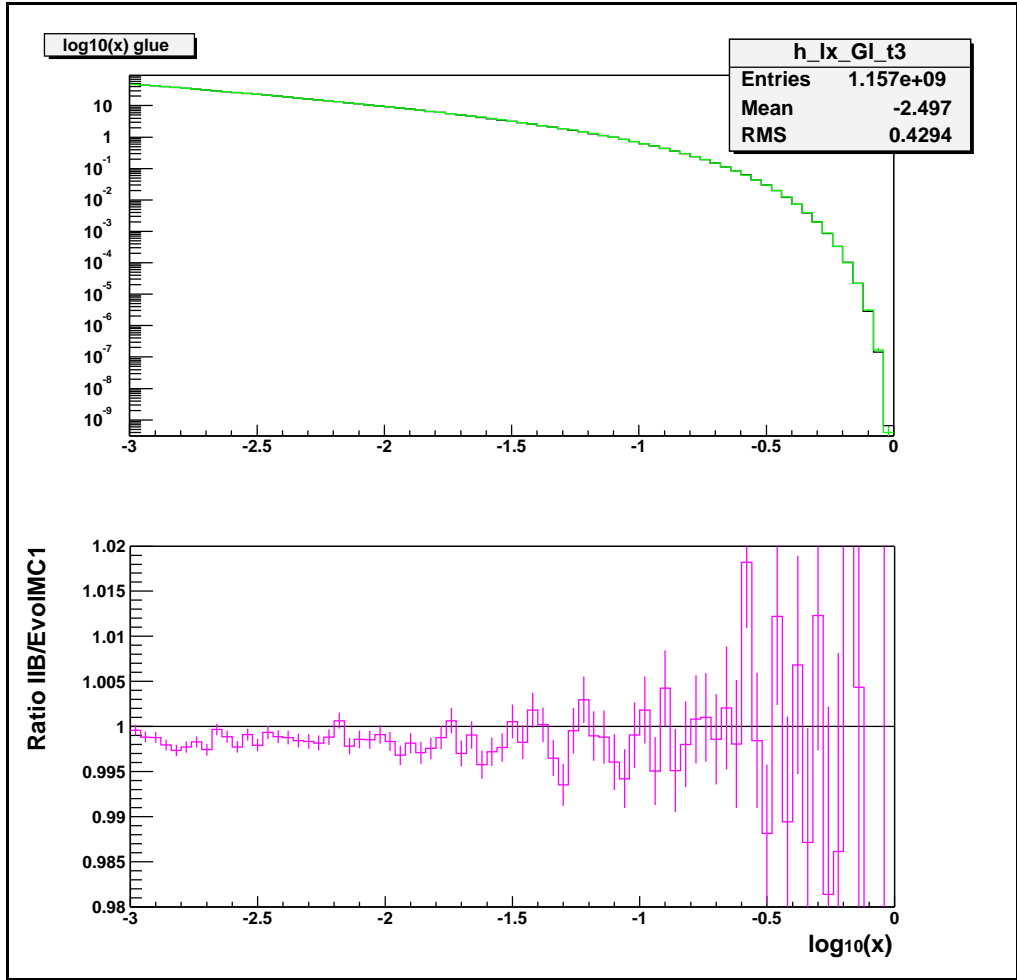


Contributions $1/z$ and $1/(1 - z)$ are combined and resummed separately.
 Worst-case scenario (pure gluon bremsstrahlung) is now prototyped and tested.

Constrained Solutions are coming

- We have found a class of solutions of the above long-standing problem
- Introductory exercise: Markovian MC $E_{\nu 0} \perp MC$ was found to agree with QCD_{num16} to within 0.2%,
[Acta Phys.Polon. B35 \(2004\) 745](#)
- Recently, 1-st prototype of the efficient constrained Markovian MC (solution IIB) prototyped.
- It agrees with the Markovian $E_{\nu 0} \perp MC$ to within 0.2%

Testing prototype IIB



Comparison of IIB solution with the Markovian MC E_{voIMC} for pure gluonstrahlung. Two solutions and the ratio (lower plot).

Agreement to within 0.2%

Short-term prospects

- More testing of IIB.
- Numerical test of solutions class I
(several solutions found, under tests)
- Implementing transitions $Q \rightarrow G$ and $G \rightarrow Q$
(at least 2 methods found)
- Adding NLL corrections
(looks rather trivial)

Long-term prospects

- Next step: Prototyping, testing and documenting the entire family of constrained MC algorithms that we see...

Long-term prospects

- Next step: Prototyping, testing and documenting the entire family of constrained MC algorithms that we see...
- Next-to-next step: looking for applications in the full scale (4-momenta) parton shower MCs. Obvious candidate processes: ISR for W/Z at LHC, DIS and ELCs.

Long-term prospects

- Next step: Prototyping, testing and documenting the entire family of constrained MC algorithms that we see...
- Next-to-next step: looking for applications in the full scale (4-momenta) parton shower MCs. Obvious candidate processes: ISR for W/Z at LHC, DIS and ELCs.
- **Bottom line:**
NEW AVENUES are opened in the construction of the **ISR PARTON SHOWER** type MCs